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The cellular ansatz: bijective combinatorics and quadratic algebra

> Xavier Viennot CNRS, LaBRI, Bordeaux

www.viennot.org

mírror websíte www.ímsc.res.ín/~víennot

Chapter 2 Quadratíc algebra, Q-tableaux and planar automata

Ch₂c

IMSc, Chennaí February 5, 2018 Xavier Viennot CNRS, LaBRI, Bordeaux <u>www.viennot.org</u>

mirror website www.imsc.res.in/~viennot

From Ch 2b

Duplication of equations in quadratic algebras

D V D $(\mathbf{U}\mathcal{P} = \mathcal{D}\mathbf{U} + \mathbf{\hat{Y}}\mathbf{\hat{X}})$ $\begin{cases} \mathbf{U} \mathbf{Y} = \mathbf{Y} \mathbf{U} \\ \mathbf{X} \mathbf{U} = \mathbf{U} \mathbf{X} \\ \mathbf{X} \mathbf{Y} = \{ \mathbf{\overline{Y}} \mathbf{\overline{X}} \}^{T} \end{cases}$ "duplication" of the commutation relations defining the algebra Q

 $U\mathcal{D} = \mathcal{D}U + \frac{1}{4} \times \frac{1}{4}$ $X_{A}Y_{A} = Y_{L}X_{L}$ $X_2 Y_2 = Y_3 X_3$ X: Y: = Yin Xin

U% = Y2U XjU = U Xj

The "RSK planar automaton" B= {Bo, B1, ... Bk} wed Bo, Aof * a = {Ao AA AA} S={0, 0} AD AO A AO AO Be Bi Aj 🖌 A: Bit A: Bi イギ」

 $V = B_0$ $X_i = B_i$ $Y_i = A_i$



DED YED $(\mathbf{U}\mathcal{P} = \mathcal{D}\mathbf{U} + \mathbf{Y}\mathbf{X})$ $\begin{vmatrix} \mathbf{U} \mathbf{Y} &= \mathbf{Y} \mathbf{U} \\ \mathbf{X} \mathbf{U} &= \mathbf{U} \mathbf{X} \\ \mathbf{X} \mathbf{Y} &= \{ \mathbf{\hat{Y}} \mathbf{\hat{X}} \}$ another "duplication" of the commutation relations of the algebra Q

・ リ フ = フ + % × $\times Y_0 = Y_1 \times$ $\begin{array}{c} X Y_{4} = Y_{2} X \\ X Y_{2} = Y_{3} X \end{array}$



 $\begin{array}{l} \times \ \gamma_{0} \ = \ \gamma_{4} \times \\ \times \ \gamma_{4} \ = \ \gamma_{2} \times \\ \times \ \gamma_{2} \ = \ \gamma_{3} \times \end{array}$ XY' Y' X



demultiplication In the PASEP algebra



 $\mathcal{D} E = E \mathcal{D} + E X_{A} + Y_{A} \mathcal{D}$ $\begin{cases} X_{A} E = E X_{2} \\ X_{i} E = E X_{i+A} \end{cases} \begin{cases} DY_{A} = Y_{2} D \\ DY_{i} = Y_{i+A} D \end{cases}$

 $X_i Y_j = Y_j X_c$









$$a_{i} = \begin{cases} 0 & \text{if no} \quad \text{in row i} \\ \neq \left[1 + number of cells - \left[-\right] \quad \text{in row i} \end{cases}$$

$$b_{j} = \begin{cases} 0 & \text{if no} \quad \text{in the } j^{\text{th}} \quad \text{column} \\ 1 + number \quad \text{f cells} \quad \left[-\right] \quad \text{in the } j^{\text{th}} \quad \text{column} \end{cases}$$

Research problem alternative Adela (T) = (P,Q) talleau size n

give a characterization of the pairs (P, Q)give a figertion between such pairs and the (n+1)! permutations of Time

Adela (T) = (P,Q)





The "Adela duality" $P(T) \leftrightarrow Q(T)$

Why Adela bijection ?



The names «Adela bijection» and «Adela duality» is in honour of my friend Adela where part of this research was done in her house in Isla Negra, Chile, inspiring place where Pablo Neruda spent many years in his house in front of the Pacific Ocean.

Isla Negra Pablo Neruda





Isla Negra Pablo Neruda

O da al nine

Nino color de día, Nino color de moche, Nino con pies de púrpura o songre de topacio, Nino, estullado hijo de la Tivra, Nino...



The 8-vertex algebra (or XYZ - algebra) (or Z - algebra)

The quadratic algebra Z.

4 generators B. A. BA 8 parameters qxy, txy {x= ,0 (BA = 900 AB + 500 A.B. B.A. = q. A.B. + t. A B $\begin{bmatrix} B, A = 9_{00} \ A B, +t_{00} \ A B \\ B A, = 9_{00} \ A, B + t_{00} \ A B.$

The Z-quadratic algebra XYZ- quadratic algebra XXZ- spin chain

Spin chains and combinatorics

A. V. Razumov, Yu. G. Stroganov

Institute for High Energy Physics 142284 Protvino, Moscow region, Russia

The XXZ quantum spin chain model with periodic boundary conditions is one of the most popular integrable models which has been investigating by the Bethe Ansatz method during the last 35 years [3]. It is described by the Hamiltonian

$$H_{XXZ} = -\sum_{j=1}^{N} \left\{ \sigma_{j}^{x} \sigma_{j+1}^{x} + \sigma_{j}^{y} \sigma_{j+1}^{y} + \Delta \sigma_{j}^{z} \sigma_{j+1}^{z} \right\}, \qquad \vec{\sigma}_{N+1} = \vec{\sigma}_{1}.$$
(1)

(XY)Z-tableaux and B.A.BA configurations (or XYZ-configurations)



 $\begin{cases} BA = 900 AB + t_{00} A.B. \\ B.A. = 900 A.B. + t_{00} A.B. \\ B.A = 900 A.B. + t_{00} A.B. \\ B.A. = 900 A.B. + t_{00} A.B. \\ BA. = 900 A.B. + t_{00} A.B. \\ BA. = 900 A.B. + t_{00} A.B. \end{cases}$



 $w \in \{B, B, A, A, F(w)\}$

W = W1 ... Wn

it step of the upper border of F(w) = } - if wi = Bor B.



bijection (s)

complete Z-talleau with diagram F(w) (w, C)B. A. BA configuration in the diagram F(W)





C = F(w) -> complete Z-talleau

w E Z B, B, A, A, J* BA

The label of ~ in T is

- if dec, try - if x & C, 9xy

 $x = \int O \quad J(\alpha) \text{ odd and } A_O$ LO $J(\alpha) \text{ even and } A_O$

50 J(x) odd and A.
5(x) even and A.

same for y with B() and BorB







......







4.16



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bijection (s)

 (\mathbf{w}, \mathbf{C}) complete Z-talleau B. A. BA configuration in the diagram F(W) with diagram F (w)

F Ferrers diagram C C F (w) - for each cell & of F - for each pair BA, BA, BA, BA, we fix a rule for the labeling of the cell by 9xy or txy (with x= or o, y= or o) according to x EC or x EC

alternating sign matrices (ASM)




t = t = 0

 $\begin{cases} BA = 900 AB + 600 A.B. \\ BA = 900 AB + 600 A.B. \\ BA = 900 AB + 600 AB \\ BA = 900 AB + 0 A.B. \\ BA = 900 AB + 0 A.B. \\ BA = 900 A.B + 0 AB. \end{cases}$

$$w = \mathbf{B}^{n} \mathbf{A}^{n}$$
 $uv = \mathbf{A}^{n}_{n} \mathbf{B}^{n}_{n}$
 $\varepsilon (u, v; w) = nb = \mathcal{A} ASM nxn$





q(A) B.A. BA configuration





q(A) B.A. BA configuration











characterisation

(i) - for any cell $x \in C$, |B(x)| and |J(x)|and same parity and

(ii) in each row and each column, the number of cells in C is odd

correlations functions in XXZ spin chains

Exact results for the σ^z two-point function of the XXZ chain at $\Delta = 1/2$

N. Kitanine¹, J. M. Maillet², N. A. Slavnov³, V. Terras⁴

14 Jun 2005

arXiv:hep-th/0506114 v1

Abstract

We propose a new multiple integral representation for the correlation function $\langle \sigma_1^z \sigma_{m+1}^z \rangle$ of the XXZ spin- $\frac{1}{2}$ Heisenberg chain in the disordered regime. We show that for $\Delta = 1/2$ the integrals can be separated and computed exactly. As an example we give the explicit results up to the lattice distance m = 8. It turns out that the answer is given as integer numbers divided by $2^{(m+1)^2}$.

¹LPTM, UMR 8089 du CNRS, Université de Cergy-Pontoise, France, kitanine@ptm.u-cergy.fr
²Laboratoire de Physique, UMR 5672 du CNRS, ENS Lyon, France, maillet@ens-lyon.fr
³Steklov Mathematical Institute, Moscow, Russia, nslavnov@mi.ras.ru

⁴LPTA, UMR 5207 du CNRS, Montpellier, France, terras@lpta.univ-montp2.fr

 e^{2z_j} , it reduces to the derivatives of order m-1 with respect to each x_j at $x_1 = \cdots = x_n = e^{\frac{\pi}{3}}$ and $x_{n+1} = \cdots = x_m = e^{-\frac{i\pi}{3}}$. If the lattice distance m is not too large, the representations (9), (11) can be successfully used to compute $\langle Q_{\kappa}(m) \rangle$ explicitly. As an example we give below the list of results for $P_m(\kappa) = 2^{m^2} \langle Q_{\kappa}(m) \rangle$ up to m = 9:

$$\begin{split} P_{1}(\kappa) &= 1 + \kappa, \\ P_{2}(\kappa) &= 2 + 12\kappa + 2\kappa^{2}, \\ P_{3}(\kappa) &= 7 + 249\kappa + 249\kappa^{2} + 7\kappa^{3}, \\ P_{4}(\kappa) &= 42 + 10004\kappa + 45444\kappa^{2} + 10004\kappa^{3} + 42\kappa^{4}, \\ P_{5}(\kappa) &= 429 + 738174\kappa + 16038613\kappa^{2} + 16038613\kappa^{3} + 738174\kappa^{4} + 429\kappa^{5}, \\ P_{6}(\kappa) &= 7436 + 96289380\kappa + 11424474588\kappa^{2} + 45677933928\kappa^{3} + 11424474588\kappa^{4} \\ &+ 96289380\kappa^{5} + 7436\kappa^{6}, \\ P_{7}(\kappa) &= 218348 + 21798199390\kappa + 15663567546585\kappa^{2} + 265789610746333\kappa^{3} \\ &+ 265789610746333\kappa^{4} + 15663567546585\kappa^{5} + 21798199390\kappa^{6} + 218348\kappa^{7}, \\ P_{8}(\kappa) &= 10850216 + 8485108350684\kappa + 39461894378292782\kappa^{2} \\ &+ 3224112384882251896\kappa^{3} + 11919578544950060460\kappa^{4} + 3224112384882251896\kappa^{5} \\ &+ 39461894378292782\kappa^{6} + 8485108350684\kappa^{7} + 10850216\kappa^{8} \\ P_{0}(\kappa) &= 911835460 + 5649499685353257\kappa + 177662495637443158524\kappa^{2} \\ &+ 77990624578576910368767\kappa^{3} + 1130757526890914223990168\kappa^{4} \end{split}$$

 e^{2z_j} , it reduces to the derivatives of order m-1 with respect to each x_i at $x_1 = \cdots = x_n = e^{\frac{\pi}{3}}$ and $x_{n+1} = \cdots = x_m = e^{-\frac{i\pi}{3}}$. If the lattice distance m is not too large, the representations (9), (11) can be successfully used to compute $\langle Q_{\kappa}(m) \rangle$ explicitly. As an example we give below the list of results for $P_m(\kappa) = 2^{m^2} \langle Q_\kappa(m) \rangle$ up to m = 9: intergers ? $P_1(\kappa) = 1 + \kappa,$ positivity ? $\mathbf{ASM} \quad P_2(\kappa) = 2 + 12\kappa + 2\beta^2,$ combinatorial interpretation $P_3(\kappa) = 7 + 249\kappa + 249\kappa^2 - 7\kappa^3,$ $P_4(\kappa) = 42 + 10004\kappa + 45444\kappa^2 + 10004\kappa^3 + 42\kappa^4,$ $P_5(\kappa) = \underbrace{429}_{738174\kappa} + 16038613\kappa^2 + 16038613\kappa^3 + 738174\kappa^4 + \underbrace{429}_{8}^5,$ $P_6(\kappa) = 7436 + 96289380\kappa + 11424474588\kappa^2 + 45677933928\kappa^3 + 11424474588\kappa^4$ $+96289380\kappa^{5}+7436\kappa^{6},$ $P_7(\kappa) = 218348 + 21798199390\kappa + 15663567546585\kappa^2 + 265789610746333\kappa^3$ (12) $+265789610746333\kappa^{4} + 15663567546585\kappa^{5} + 21798199390\kappa^{6} + 218348\kappa^{7},$ $P_8(\kappa) = 10850216 + 8485108350684\kappa + 39461894378292782\kappa^2$ $+ 3224112384882251896\kappa^3 + 11919578544950060460\kappa^4 + 3224112384882251896\kappa^5$ $+39461894378292782\kappa^6+8485108350684\kappa^7+10850216\kappa^8$ $P_9(\kappa) = 911835460 + 5649499685353257\kappa + 177662495637443158524\kappa^2$ $+77990624578576910368767\kappa^3+1130757526890914223990168\kappa^4$

The number of B.A. BA configurations in the grid [n]x[n] is





rhombus tilings







coding of the edges for a tiling of the triangular lattice



rewriting rules for tilings of the triangular lattice



BA=AB+A.B.





serviting rules for tilings of the triangular lattice



rewriting rules for tilings of the triangular lattice

$$BA=AB+A_B$$

 $B_A=AB$
 $B_A=AB$

 $\int t_{00} = t_{00} = 0$ (ASM) $\int q_{00} = 0$

Rhombus tilings

 $\begin{cases} BA = 900 AB + 600 A.B. \\ B.A. = 0 A.B. + 600 A.B. \\ B.A = 900 A.B. + 600 A.B. \\ B.A. = 900 AB. + 0 A.B. \\ BA. = 900 A.B. + 0 A.B. \\ BA. = 900 A.B. + 0 AB. \end{cases}$







w = BBABAC A^c A^b B^a B^c c(u,v;w)



plane partitions

64331 4221









The number of plane partitions in a box (a, b, c)



Proof of MacMahon formula, see: BJC 1, Ch 5a, p105



in a box a x b x c

151 6a

15/56

15RSC

i+j+k-1 i+j+k-2

= c(u,v;w)

$$BA = AB + A_B,$$
$$B_A = AB$$
$$B_A = A_B,$$
$$B_A = AB,$$



dímers tiling on a square lattice

a tiling (of a rectangle) on the lettice





for tilings (square lattice)















BA = AB



BA = AB + AB



BA = AB + ABBA = AB of the Z-algebra BA. = AB A = 0Correction to the video: 900 A B 16 Formula for the number of tilings B. A. = AB of an mxn rectangle, B, A A.B see BJC 1, Ch 5b, p66-67 (without proof) 90. A. B B AB.

Aztec tilings










(Aztec lattice)





Aztec tilings too = too = 0 t. = 2

An (2) enumeration of ASM according to the number of (-1)

A (2) n(n-1)/2









geometric interpretation of XYZ-tableaux:

6 and 8 vertex models



The 8- vertex model

 $BA = q_{00} AB + t_{00} A.B.$ $\begin{array}{c} B_{\bullet}A_{\bullet} = q_{\bullet\bullet} & A_{\bullet}B_{\bullet} + t_{\bullet\bullet} & A & B \\ B_{\bullet}A = q_{\bullet\bullet} & A & B_{\bullet} + t_{\bullet\bullet} & A & B \\ B_{\bullet}A = q_{\bullet\bullet} & A & B_{\bullet} + t_{\bullet\bullet} & A & B \\ B_{\bullet}A_{\bullet} = q_{\bullet\bullet} & A_{\bullet}B + t_{\bullet\bullet} & A & B_{\bullet} \end{array}$



The 6-vertex model

$$\begin{cases} B A = 9 & A B + t & A B \\ B A = 9 & A B + t & A B \\ B A = 9 & A B + t & A B \\ B A = 9 & A B + O A B \\ B A = 9 & A B + O A B \\ B A = 9 & A B + O A B \\ B A = 9 & A B + O A B \\ \end{bmatrix}$$







The 6-vertex model



the ice model

H & H & H O H O H O H H-S HOH



bijection

configurations of the 6-vertex model









